

Physics Beyond the Standard Model —

Exercise Sheet 5

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1. **Spinor and Superfield Calculations** Carry out as many of the calculations below as you like! Focus in particular on parts (c) and (e).

(a) Rewrite the QED-Lagrangian

$$\mathcal{L}_{\text{QED}} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + \bar{\Psi}(i\not{D} - m)\Psi \quad (1)$$

in terms of two-component Weyl-spinors

$$\Psi = \begin{pmatrix} \Psi_{L\alpha} \\ \bar{\Psi}_{\dot{R}\dot{\alpha}} \end{pmatrix} \quad \text{and} \quad \bar{\Psi} = \begin{pmatrix} \Psi_R^\alpha & \bar{\Psi}_{L\dot{\alpha}} \end{pmatrix}.$$

(b) Using

$$\epsilon_{\alpha\beta} = -\epsilon_{\beta\alpha} \quad , \quad \epsilon_{12} = 1 \quad (2)$$

$$\epsilon^{\alpha\beta} \epsilon_{\beta\gamma} = \delta^\alpha_\gamma \quad , \quad \epsilon^{\alpha\beta} = \epsilon_{\beta\alpha} \quad (3)$$

$$\psi^\alpha = \epsilon^{\alpha\beta} \psi_\beta \quad (4)$$

$$\psi\chi = \psi^\alpha \chi_\alpha = -\chi_\alpha \psi^\alpha = \chi^\alpha \psi_\alpha = \chi\psi \quad (5)$$

show the following identities:

$$\theta^\alpha \theta_\beta = \frac{1}{2} \delta^\alpha_\beta \theta\theta \quad (6)$$

$$\bar{\theta}_{\dot{\alpha}} \bar{\theta}^{\dot{\beta}} = \frac{1}{2} \delta_{\dot{\alpha}}^{\dot{\beta}} \bar{\theta}\bar{\theta} \quad (7)$$

$$\partial_\alpha \theta\theta = 2\theta_\alpha \quad (8)$$

$$\bar{\partial}^{\dot{\alpha}} \bar{\theta}\bar{\theta} = 2\bar{\theta}^{\dot{\alpha}} \quad (9)$$

$$\theta\sigma^\mu \bar{\psi} = -\bar{\psi}\bar{\sigma}^\mu \theta \quad (10)$$

$$\psi\theta\theta_\alpha = -\frac{1}{2}\theta\theta\psi_\alpha \quad (11)$$

$$\bar{\psi}\bar{\theta}\bar{\theta}^{\dot{\alpha}} = -\frac{1}{2}\bar{\theta}\bar{\theta}\bar{\psi}^{\dot{\alpha}}. \quad (12)$$

(c) Show $\{Q_\alpha, \bar{Q}_{\dot{\beta}}\} = 2\sigma_{\alpha\dot{\beta}}^\mu P_\mu$ using

$$P^\mu = i\partial^\mu \quad (13)$$

$$Q_\alpha = i(\partial_\alpha + i_\alpha(\sigma^\mu \bar{\theta})\partial_\mu) \quad (14)$$

$$\bar{Q}_{\dot{\alpha}} = i(-\bar{\partial}_{\dot{\alpha}} - i(\theta\sigma^\mu)_{\dot{\alpha}}\partial_\mu). \quad (15)$$

(d) Show that the following Lagrangian $\mathcal{L} = \bar{\psi}i\bar{\sigma}^\mu \partial_\mu \psi + |\partial_\mu \phi|^2 + |F|^2$ is invariant (up to total derivatives) under the supersymmetry transformation

$$\delta\phi = \sqrt{2}\epsilon\psi, \quad \delta\psi^\alpha = \sqrt{2}\epsilon^\alpha F - \sqrt{2}i(\bar{\epsilon}\bar{\sigma}^\mu)^\alpha \partial_\mu \phi, \quad \delta F = -i\sqrt{2}\bar{\epsilon}\bar{\sigma}^\mu \partial_\mu \psi$$

- (e) Chiral superfields are superfields with a particularly simple structure. They are used to describe e.g. the electron-selectron or Higgs-Higgsino multiplets. Two chiral superfields $\Phi_{L,R}$ are defined as follows:

$$\Phi_{L,R}(x, \theta, \bar{\theta}) = \exp(-i\theta\sigma^\mu\bar{\theta}\partial_\mu)\phi_{L,R}(x, \theta) \quad (16)$$

with

$$\phi_{L,R}(x, \theta) = A(x) + \sqrt{2}\theta\psi(x) + \theta\theta F(x). \quad (17)$$

- i. Compute the $\theta\theta$ -part of the product $\Phi_L\Phi_R$, which is defined to be

$$\int d^2\theta \Phi_L\Phi_R|_{\bar{\theta}=0}$$

This is the electron mass term in SUSY QED.

- ii. Compute Φ_L by evaluating the exponential as a power series.
 iii. Calculate the $\theta\theta\bar{\theta}\bar{\theta}$ -part of the product $\Phi_L^\dagger\Phi_L$, defined to be

$$\int d^4\theta \Phi_L^\dagger\Phi_L \equiv \int d^2\theta d^2\bar{\theta} \Phi_L^\dagger\Phi_L.$$

Show that the result is equal to the Lagrangian given in part (d).

2. Type X Yukawa interaction in the 2HDM

In the type X Yukawa interaction we impose the Z_2 transformation: $\Phi_1 \rightarrow -\Phi_1$, $e_{jR}^0 \rightarrow -e_{jR}^0$, and $\mathcal{S} \rightarrow \mathcal{S}$ for all the other fields. The Z_2 -invariant Yukawa interaction term can be written as

$$-\mathcal{L}_Y^X = \bar{Q}_i^0 \Delta_{ij}^u \Phi_2^c u_{jR}^0 + \bar{Q}_i^0 \Delta_{ij}^d \Phi_2 d_{jR}^0 + \bar{L}_i^0 \Gamma_{ij}^d \Phi_1 e_{jR}^0. \quad (18)$$

- (a) Referring to the lecture notes, diagonalize the mass matrices and calculate the Yukawa interaction couplings of the fermions to the neutral and charged Higgs fields.
 (b) As a phenomenological example we show here a Barr-Zee type Feynman diagram contributing to the muon anomalous magnetic moment. For which fermion f in the inner loop is the Feynman diagram enhanced by $\tan^2\beta$? (Note that $\tan\beta$ can be realistically as large as ~ 60 .)

