

ATQ 2022/23 - sheet 6

(a) 1D continuity equation $e \frac{\partial \rho(x,t)}{\partial t} = -\partial_x J(x,t)$

With $\rho(x,t) \approx \rho_R(x,t) + \rho_L(x,t) = \frac{1}{\sqrt{\pi}} \partial_x \phi(x,t)$, we get:

$$-\partial_x \left(-\frac{e}{\sqrt{\pi}} \partial_t \phi(x,t) \right) = -\partial_x J(x,t)$$

$$\leadsto J(x,t) = -\frac{e}{\sqrt{\pi}} \partial_t \phi(x,t) \quad (\text{up to integration constant})$$

{no zero-order
equilibrium current}

✓

(b) Fourier transform $\langle J(x) \rangle(t) = \langle J(x) \rangle_0(t) + \int dx' \int dt' G_{JJ}^R(x-x', t-t') A(x', t')$

with $G_{JJ}^R(x-x', t-t') = -i \theta(t-t') \langle [J(x,t), J(x', t')] \rangle_0 = -\frac{ie^2}{\pi} \theta(t-t') \langle [\partial_t \phi(x,t), \partial_{t'} \phi(x', t')] \rangle_0$

(a)

$$\langle J(x) \rangle(\omega) = \int dt e^{i\omega t} \langle J(x) \rangle(t) = \int dt e^{i\omega t} \int dx' \int dt' G_{JJ}^R(x-x', t-t') A(x', t') =$$

$$\begin{aligned}
 &= \int dx' G_{jj}^R(x-x', \omega) A(\omega) = \int dx' G_{jj}^R(x-x', \omega) \frac{-i E(\omega)}{\omega} \checkmark \\
 &\uparrow \\
 &\left[(G * A)(t) \text{ (time-convolution)} \right] \\
 &\leadsto (G * A)(\omega) = G(\omega) A(\omega) \\
 &\uparrow \\
 &\left[\text{convolution theorem} \right]
 \end{aligned}$$

$$(c) \quad Z = \text{tr}(e^{-\beta H}) = \int D\phi D\theta e^{-\int_0^\beta d\tau \int dx [-i \partial_x \theta \partial_\tau \phi + H_0(\phi, \theta)]}$$

$$H_0(\phi, \theta) = \frac{\mu}{2} \left[\kappa (\partial_x \theta)^2 + \frac{1}{\kappa} (\partial_x \phi)^2 \right]$$

Think of $\int D\phi$ as $\lim_{N_1, N_2 \rightarrow \infty} \int \left(\prod_{j_1=0}^{N_1-1} \prod_{j_2=0}^{N_2-1} d\phi(x_{j_1}, \tau_{j_2}) \right)$, similar for $D\theta$

↳ (See Altland and Simons Chapter 3 + 4 for a pedagogical introduction)

$$\left[\text{In reciprocal space } \int D\phi \rightarrow \int \prod_{q, \omega_n} d\phi_{q, i\omega_n} \right]$$

$\uparrow$
 discrete for finite L

- Eliminate θ by Gaussian integration or using the

Heisenberg equations of motion $\partial_\tau \phi(x, \tau) = [H, \phi] = \int dx' \frac{\partial H}{\partial (\partial_{x'} \theta)} [\partial_{x'} \theta(x', \tau), \phi(x, \tau)]$

$= -i u K \partial_x \theta(x, \tau)$, i.e. $-i \partial_x \theta = \frac{1}{u K} \partial_\tau \phi(x, \tau)$, and thus

$$\left[\phi(x), \partial_{x'} \theta(x') \right] = i \delta(x - x') \quad -i \partial_x \theta \partial_\tau \phi - H_0(\phi, \theta) = \frac{u}{2K} \left(\frac{1}{u^2} (\partial_\tau \phi)^2 + (\partial_x \phi)^2 \right) \quad \checkmark^{(11)}$$

• To reciprocal space: $\phi(x, \tau) = \sum_{q, i\omega_n} e^{i(qx - \omega_n \tau)} \phi_{q, i\omega_n}$

$$\leadsto \int d\tau \int dx \frac{u}{2K} \left(\frac{1}{u^2} (\partial_\tau \phi)^2 + (\partial_x \phi)^2 \right) = \beta L \frac{u}{2K} \sum_{q, i\omega_n} \left(q^2 + \frac{\omega_n^2}{u^2} \right) \phi_{-q, -i\omega_n} \phi_{q, i\omega_n}$$

$$\left[\int_0^\beta d\tau e^{i(\omega_n + \omega_m)\tau} = \beta \delta_{\omega_n, -\omega_m} ; \int_0^L dx e^{i(q+q')x} = L \delta_{q, -q'} \right]$$

• Compute $\langle \phi_{-q, -i\omega_n} \phi_{q', i\omega_{n'}} \rangle_0 = \frac{1}{Z} \int \left(\prod_{\vec{q}, i\omega_n} d\phi_{\vec{q}, i\omega_n} \right) e^{-\frac{\beta \mathcal{L}}{2K} \sum_{\vec{q}, i\omega_n} \left(\vec{q}^2 + \frac{\omega_n^2}{u^2} \right) \phi_{\vec{q}, i\omega_n}^2} = \dots$

Gaussian integration rule:

$$\int \left(\prod_j d\vec{u}_j \right) e^{-\frac{1}{2} \vec{u}^T A \vec{u}} = \frac{1}{\sqrt{\det A}}$$

$$\int \left(\prod_j d\vec{u}_j \right) e^{-\frac{1}{2} \vec{u}^T A \vec{u}} = \left(A^{-1} \right)_{x,y}$$

only $\phi_{q>0, \omega_n>0}$ independent
 complex variables, since $\phi(x, \tau) \in \mathbb{R}$
 Here, $u_x^* \rightarrow \phi_{-q, -i\omega_n}$, $u_y \rightarrow \phi_{q, i\omega_n}$
 $A_{xy} \rightarrow \delta_{q, q'} \delta_{n, n'} \frac{\beta u}{K} \left(q^2 + \frac{\omega_n^2}{u^2} \right)$

$$= \delta_{q, q'} \delta_{n, n'} \frac{K}{\beta u} \frac{1}{q^2 + \frac{\omega_n^2}{u^2}} \quad (*)$$

• With that compute current-current correlator:

$$\langle T_M (\partial_\tau \phi(x, \tau) \partial_{\tau'} \phi(x', \tau')) \rangle_{T=0} = \left\langle T_M \left(\sum_{q, i\omega_n} (-i\omega_n) \phi_{q, i\omega_n} e^{i(qx - \omega_n \tau)} \sum_{q', i\omega_{n'}} (-i\omega_{n'}) e^{iq'x'} \phi_{q', i\omega_{n'}} \right) \right\rangle_0$$

$$\begin{aligned}
 & \xrightarrow{(\#)} \sum_{q, \omega_n} (\omega_n^2) e^{i[q(x-x') - \omega_n \tau]} \quad \frac{K}{\beta L u} \frac{1}{q^2 + \frac{\omega_n^2}{u^2}}
 \end{aligned}$$

$\left[\begin{array}{l} n = -n' \\ q' = -q \end{array} \right]$

$$\leadsto G_{yy}^M(x-x', i\omega_n) = -\frac{e^2}{\pi} \int d\tau e^{i\omega_n \tau} \langle T_M(\partial_\tau \phi(x, \tau) \partial_{\tau'} \phi(x', \tau'))|_{\tau'=0} \rangle =$$

$$= -\frac{e^2}{\pi} \int_q (\omega_n)^2 e^{iq(x-x')} \frac{K}{L u} \frac{1}{q^2 + \frac{\omega_n^2}{u^2}} = -\frac{e^2}{\pi} \frac{K \omega_n^2}{u} \int \frac{dq}{2\pi} e^{iq(x-x')} \frac{1}{q^2 + \frac{\omega_n^2}{u^2}}$$

$$\left[\sum_q \rightarrow \frac{L}{2\pi} \int dq \right]$$

$$\xrightarrow{\uparrow} = -\frac{e^2 K \omega_n}{2\pi} e^{-|x-x'| \frac{\omega_n}{u}} \quad \checkmark$$

$$\left[\int \frac{dq}{2\pi} e^{iq(x-x')} \frac{1}{q^2 + a^2} = \frac{1}{2a} e^{-a|x-x'|} \right]$$

(d) Analytic continuation:

$$G_{jj}^R(x-x', \omega) = G_{jj}^M(x-x', i\omega_n \rightarrow \omega + i0^+) \stackrel{(c)}{\downarrow} = - \frac{e^2 K (\omega + i0^+)}{i 2\pi} e^{i \frac{\omega + i0^+}{u} |x-x'|}$$

Plug back into Kubo formula:

$$\langle J(x) \rangle(\omega) = \int dx' E(x', \omega) \cancel{\frac{i}{\omega}} \frac{e^2 K \cancel{\omega}}{2\pi \cancel{\omega}} e^{i \frac{\omega}{u} |x-x'|} \stackrel{\substack{\text{[DC-limit } \omega \rightarrow 0] \\ \text{[uniform } E]}}{\downarrow} =$$

$$= \underbrace{E \frac{e^2 K}{2\pi} L}_{= EG} ; \quad G = \frac{G}{L} = \frac{e^2}{2\pi} K \quad \checkmark$$