

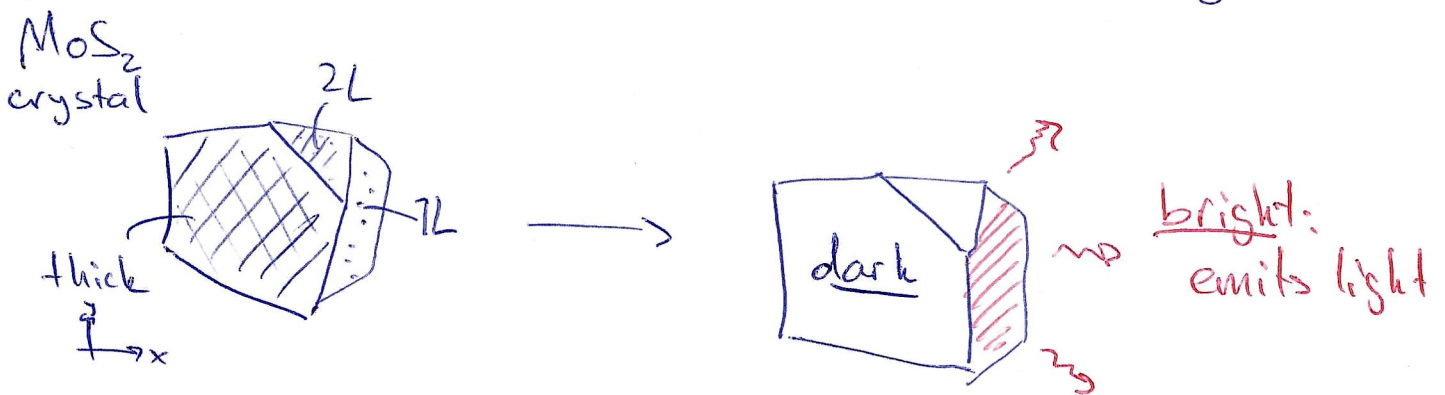
4. Atomically-thin semiconductors

- ↳ materials for opto-electronics, photonics, valleytronics* (valley-index as information carrier), and twistronics* (rotation angle to tune material properties)
- ↳ strong electronic interactions & light-matter coupling (=> fundamental solid state physics)

main representatives: transition metal dichalcogenides (TMDCs: $\text{MoS}_2, \text{WSe}_2, \dots$)

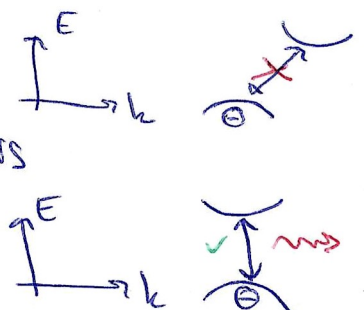
① Initial observation:

|| Only monolayers are efficient light emitters



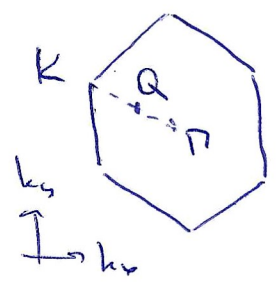
↳ reason: → Bulk, 4L, 3L, 2L... are indirect semiconductors

→ 1L are direct semiconductors



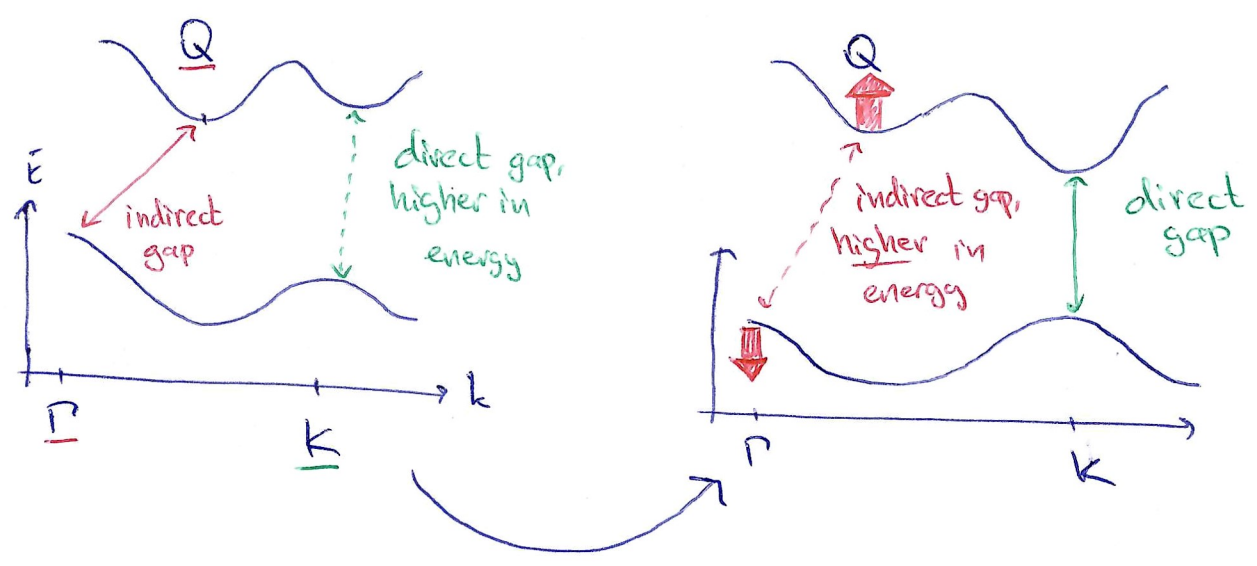
② Indirect-to-direct transition

Q: "Why does it happen?"



Bulk, ..., 2L

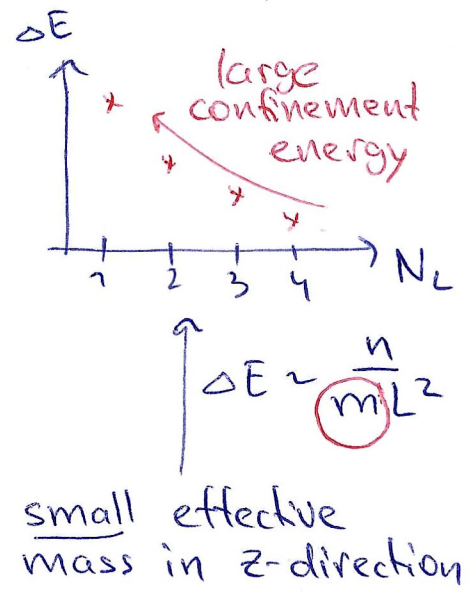
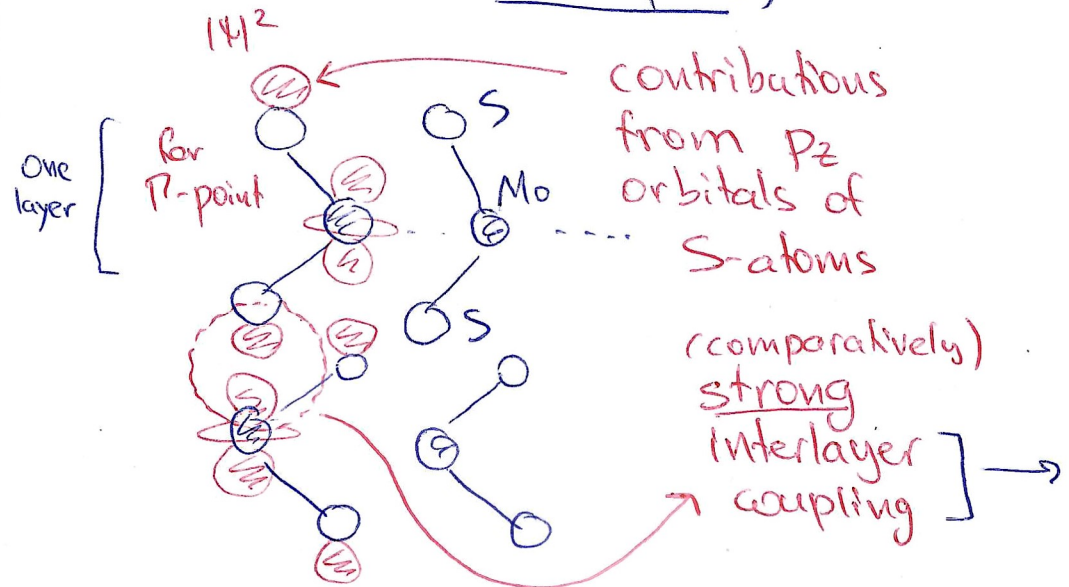
1L



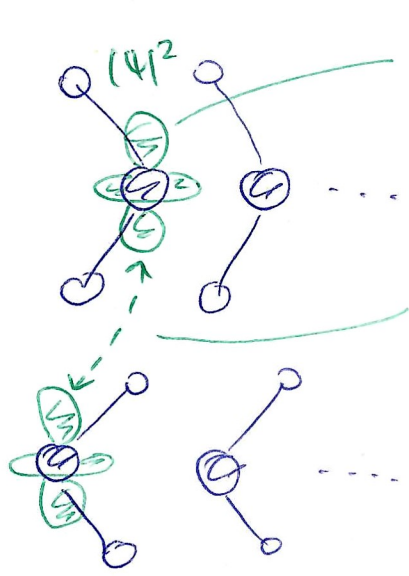
- indirect gap increases
- direct gap remains unaffected

why?

↳ indirect gap: electron states at Γ and Q points: (wavefct. in real space)

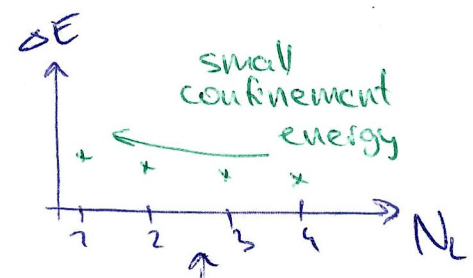


↳ **direct gap** : electron states at K points



d-orbitals of Mo-atoms

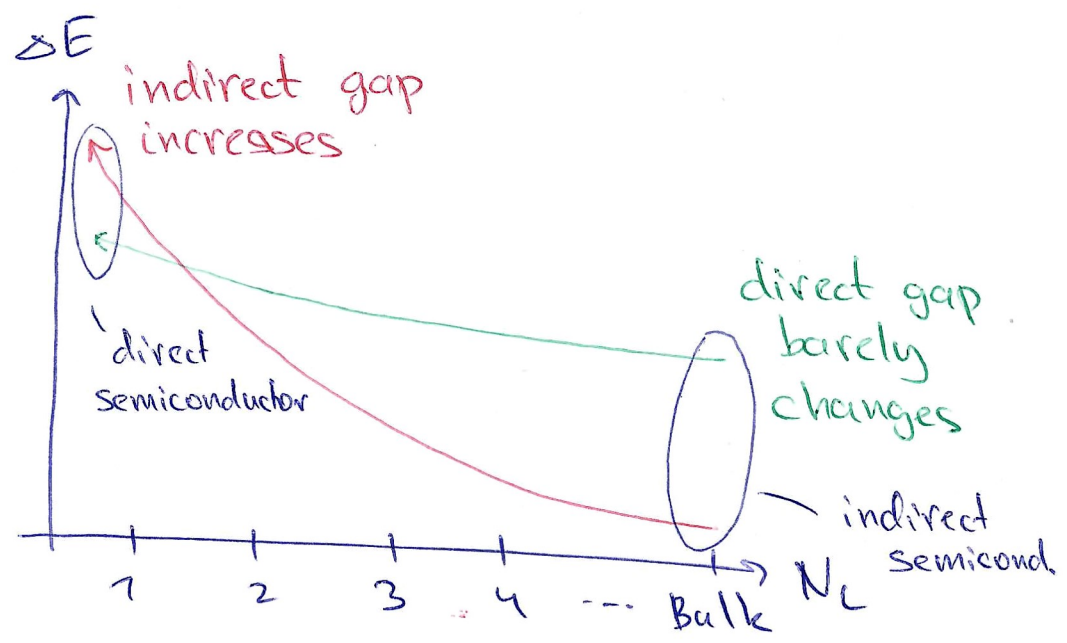
no spatial overlap, weak interlayer coupling



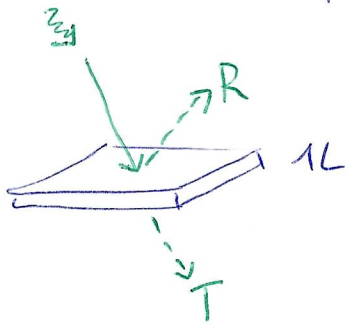
$\Delta E \sim \frac{\hbar^2}{mL^2}$

large effective mass in z-direction

→
to summarize



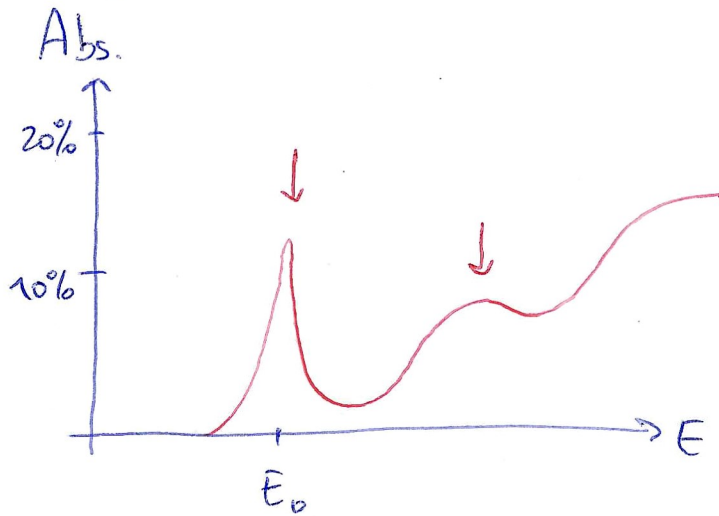
③ Optical absorption of TMDC monolayers



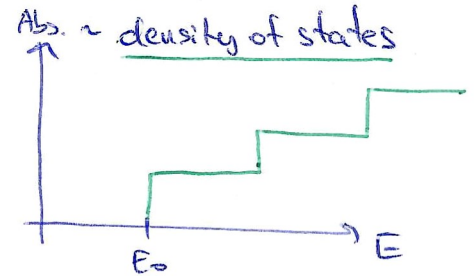
absorption

$$A = 1 - R - T$$

→ typical spectrum at room temperature



vs. expectation from non-interacting 2D electrons:



→ interactions between electrons (and holes) are strong:

Coulomb forces:

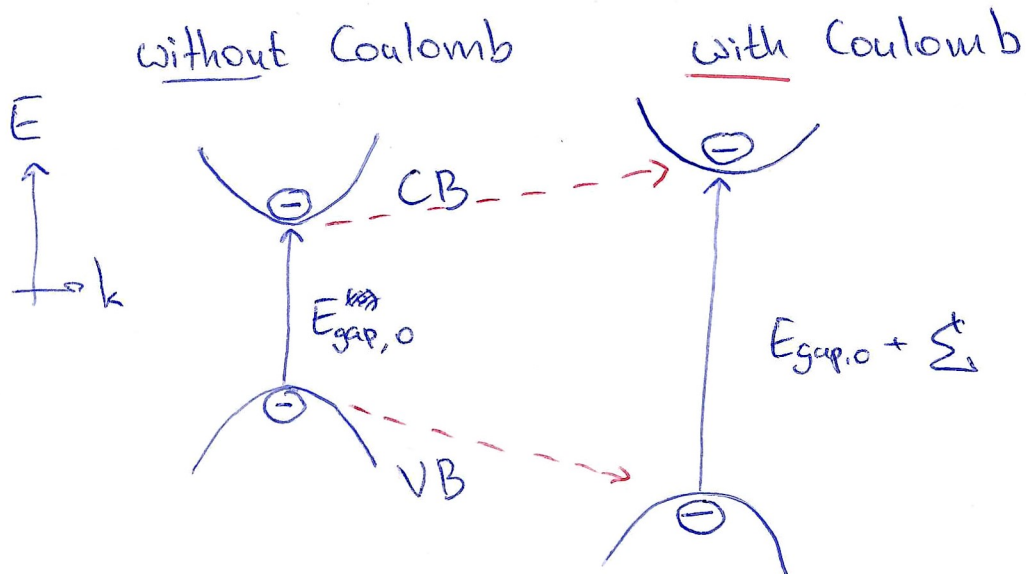
Diagram illustrating Coulomb interactions: two electrons (⊖) repel each other (← ⊖ ⊖ →), and an electron (⊖) and a hole (⊕) attract each other (⊖ → ← ⊕). A distance 'r' is indicated between the electron and hole.

$$\left[\vec{F} \sim \frac{e^2}{r^2} \quad (\text{in 3D}) \right]$$

→ Strong Coulomb interactions lead to renormalization (=change) of energies & formation of new quasiparticles

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Bandgap renormalization

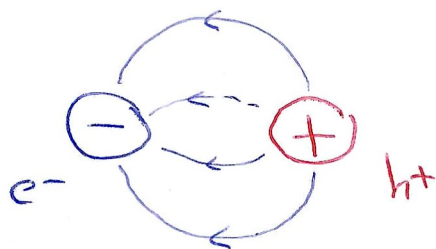


→ energy required to excite an e^- from valence to conduction band increases due to e^-e^- -interactions (contribution Σ , typically $\Sigma < E_{\text{gap},0}$)

|| → effective size of the bandgap depends on the strength of Coulomb interaction

New quasiparticles: Excitons

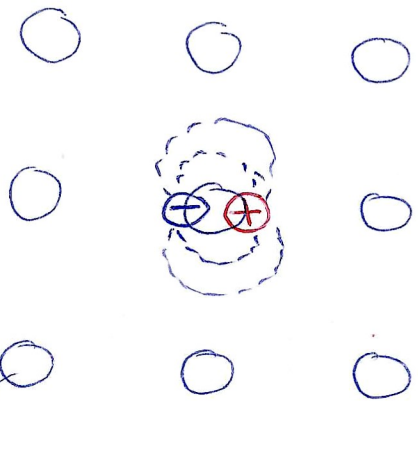
Excitons := Coulomb-bound electron-hole pairs (from attractive interaction)



→ often determine optical properties of semiconductors

- generally two main types of excitons:

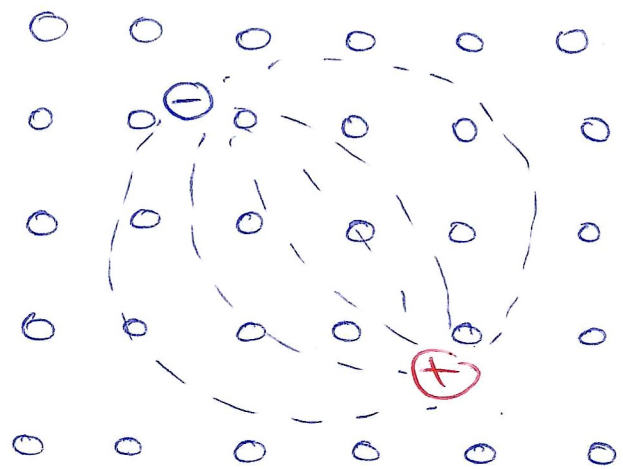
Frenkel exciton



molecule/unit cell
in a crystal

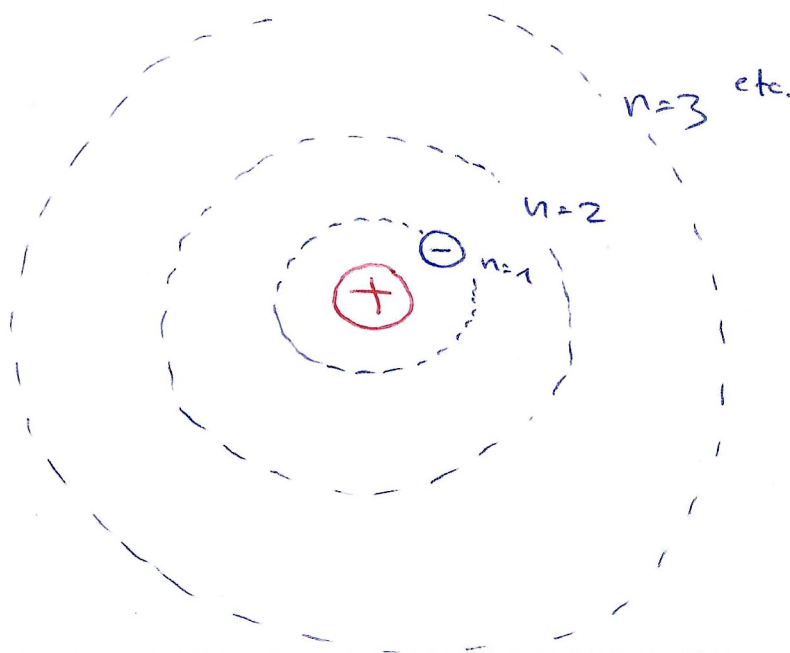
- localized on one unit cell

Wannier-Mott exciton



- delocalized over many unit cells

→ in 2D materials excitons can share both Frenkel & Wannier properties, yet can be often described accurately in Wannier picture:



≅ hydrogenic model =

similar to proton+electron problem

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 → However in contrast to H-atom, exciton are composed of e^- and h^+ quasiparticles with effective masses m^* that are located in a solid state environment with increased dielectric screening ϵ

→ Binding energies* in 2D can be approximated as:

* how strong e^- and h^+ are bound

$$E_{\text{bind}} = \frac{R_y^{-13.6\text{eV}}}{(n - \frac{1}{2})^2} \times \frac{\mu}{\epsilon^2}$$

• $\frac{1}{(n - \frac{1}{2})^2}$ from 2D dimensionality

• $\mu = \left(\frac{1}{m_e} + \frac{1}{m_h}\right)^{-1}$ reduced effective mass

• ϵ : dielectric screening constant of the environment

* typical values for quantum wells (e.g. GaAs)

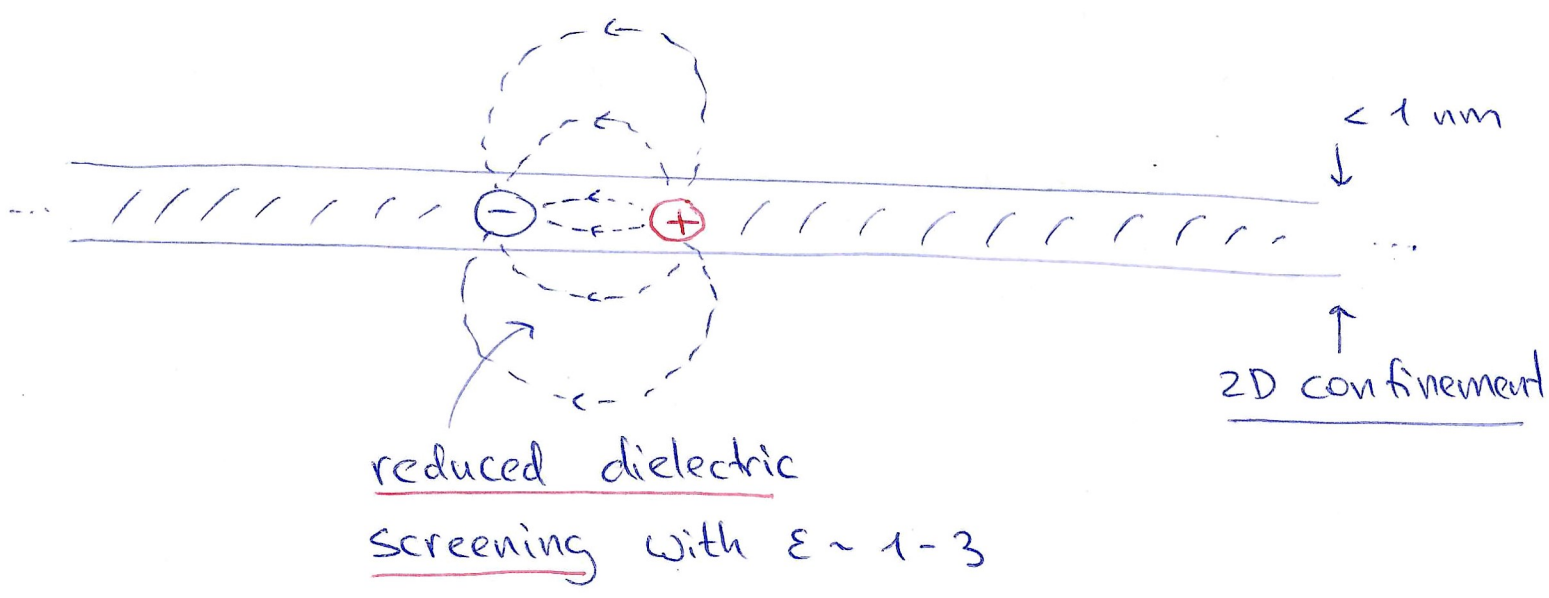
$$\mu \sim 0.1 m_0$$

$$\epsilon \sim 10$$

$$\rightarrow E_{\text{bind}} \sim 10^1 \text{ of meV} \ll E_{\text{gap}}, \leq k_B T$$

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→ excitons in monolayer materials

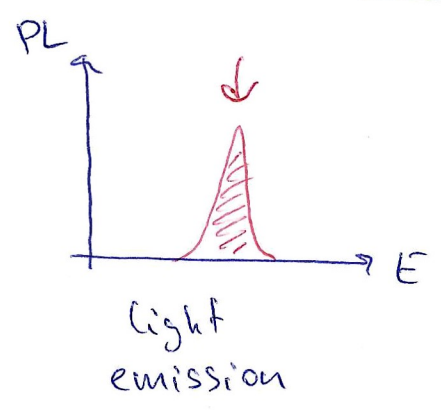
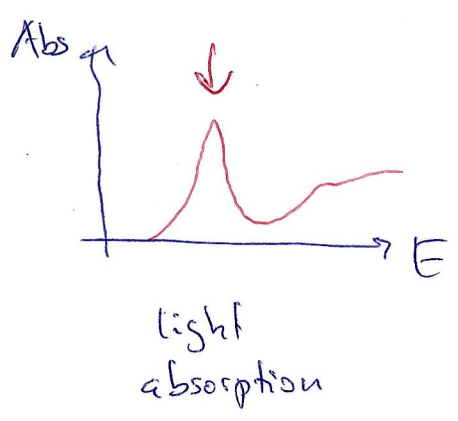


→ binding energies are very large,
0.2 - 0.5 eV ($\gg k_B T$)

④ Measuring Coulomb effects

option 1: Combine optics & electron transport

↳ Exciton ground state from optical spectra

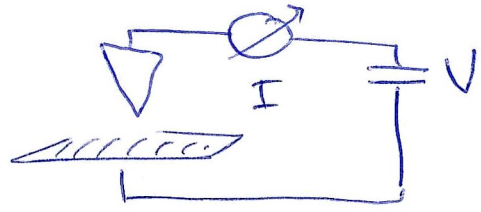


✓

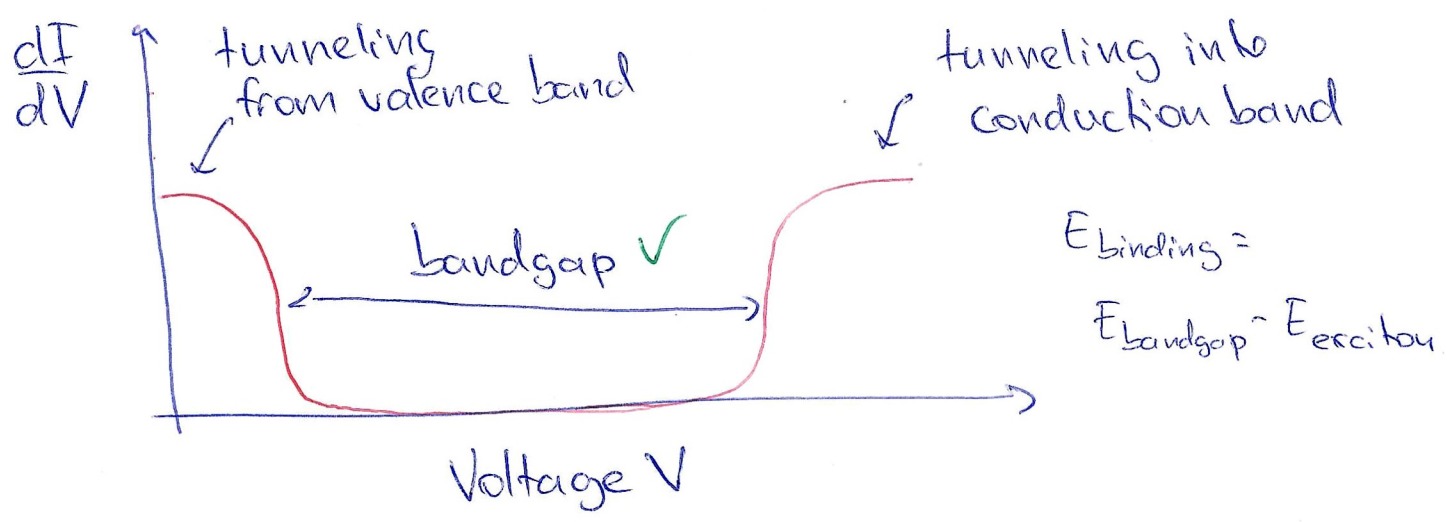
g

→ electronic bandap from scanning tunneling spectroscopy* ✓

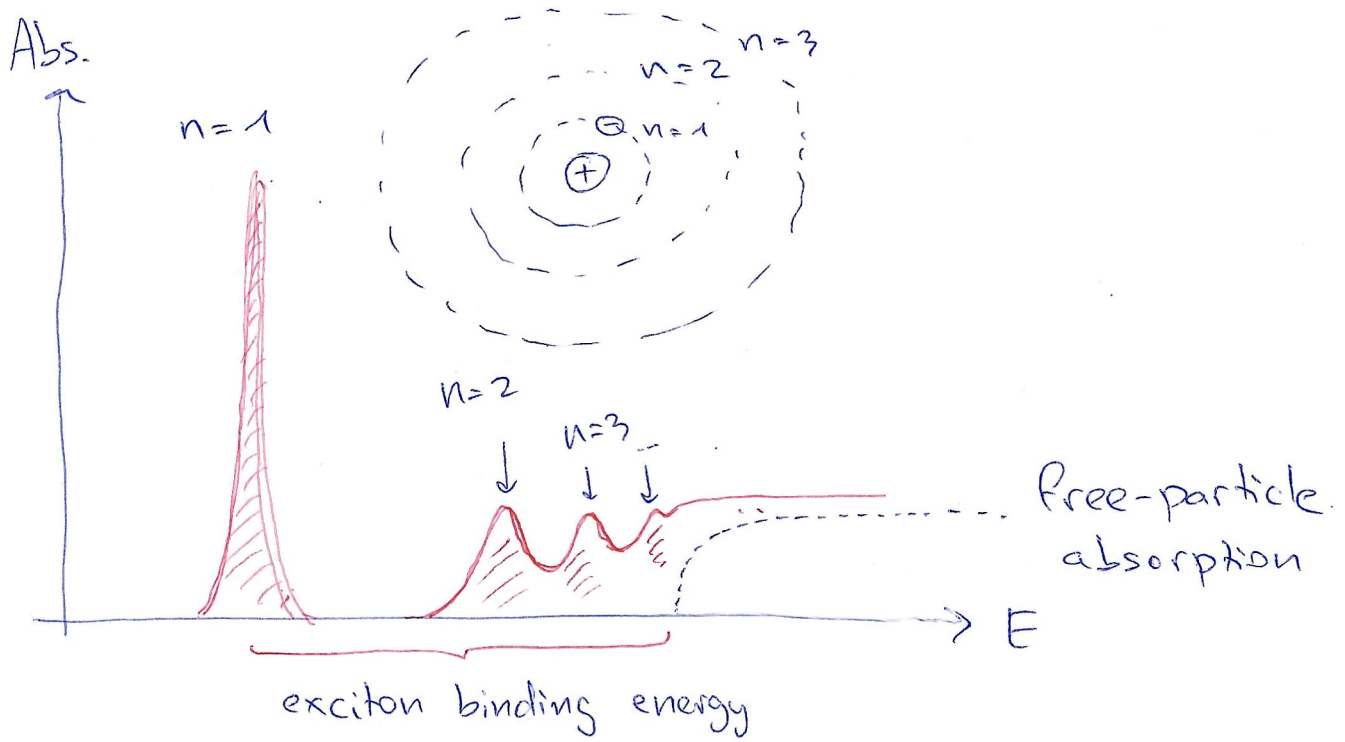
tunes the possibility to tunnel through possible transitions into from occupied into unoccupied states



* := $I(V)$ one one position



Option 2: Measure excited exciton states

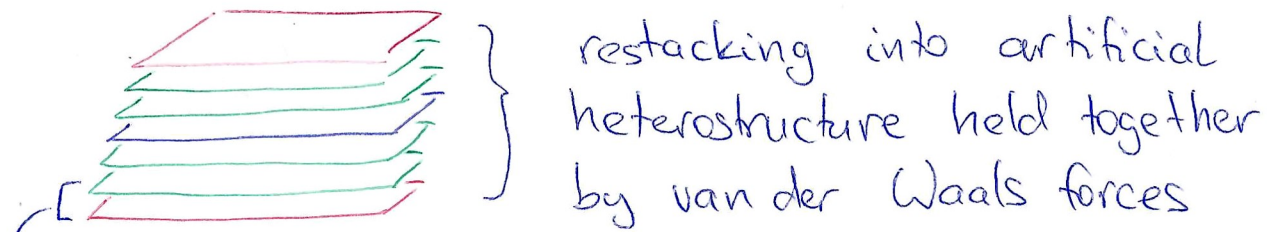


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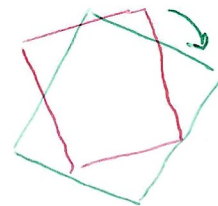
④ Heterostructures

→ artificial combinations of different materials (including layered systems) has a rich history of research (e.g. quantum wells, organic blends, etc...)

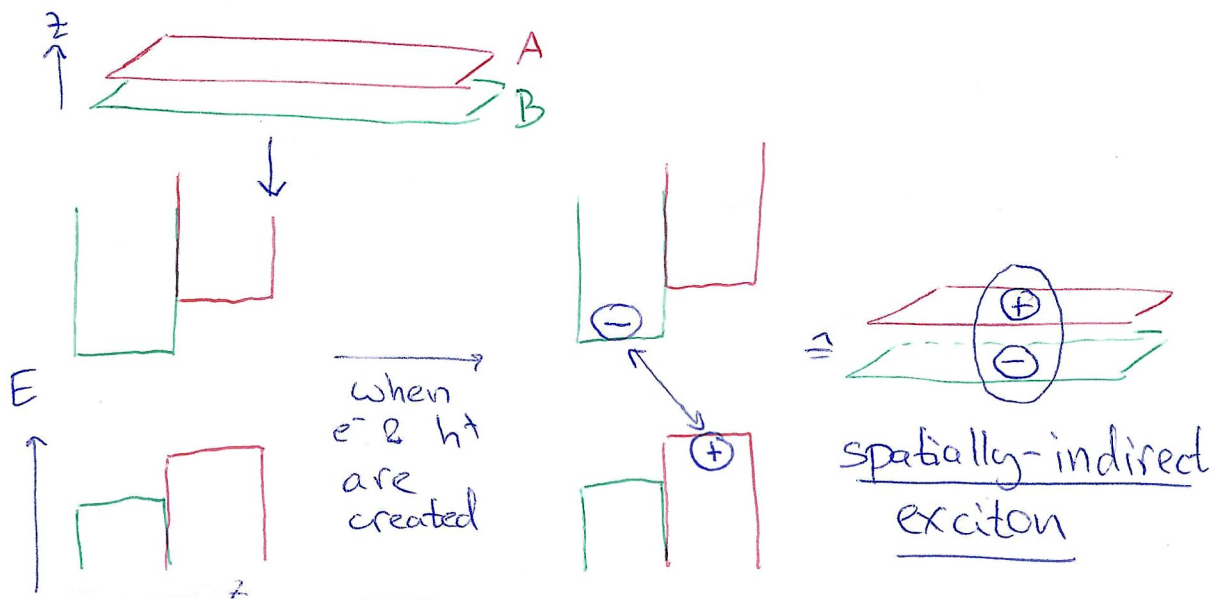
→ Van der Waals crystals offer an excellent platform to combine different materials due to relaxed lattice matching requirements



interfaces can be atomically sharp + layers can be twisted



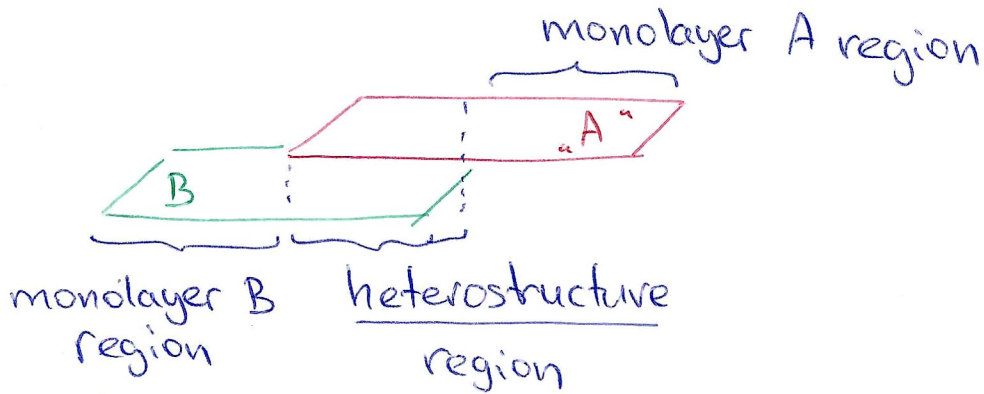
→ e.g. combining two semiconducting layers:



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↳ indirect excitons are:

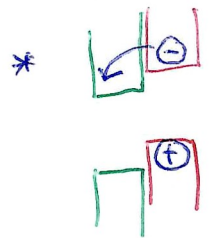
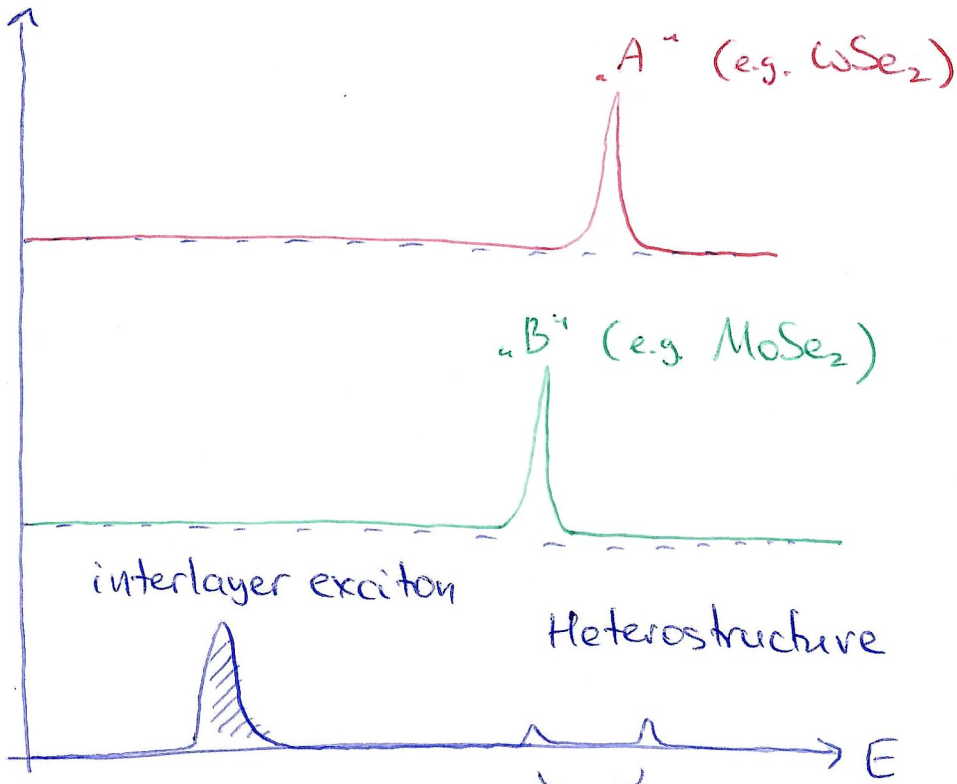
- (often) long-lived
- externally highly tunable
- can store spin- or valley "information" for a long time

Q: "How to detect interlayer excitations?"



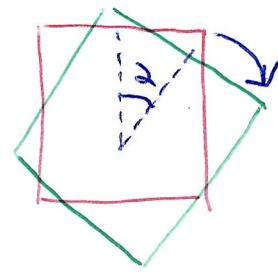
(light emission)

PL



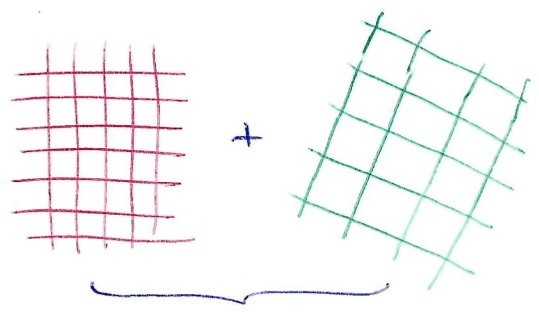
reduced intra-layer PL due to efficient charge transfer*

→ electronic, optical, and vibrational properties can be tuned by the relative twist angle between the layers

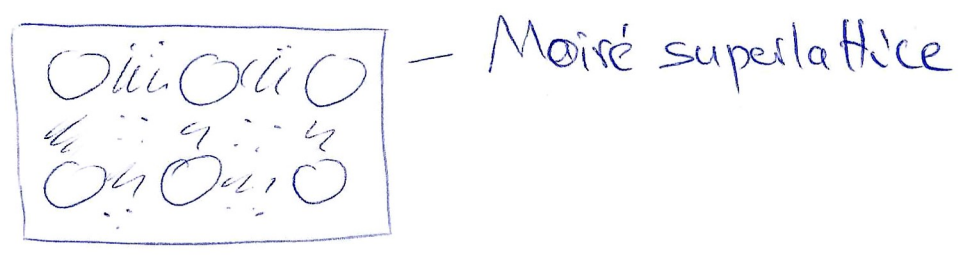


→ important phenomenon: Moiré effect

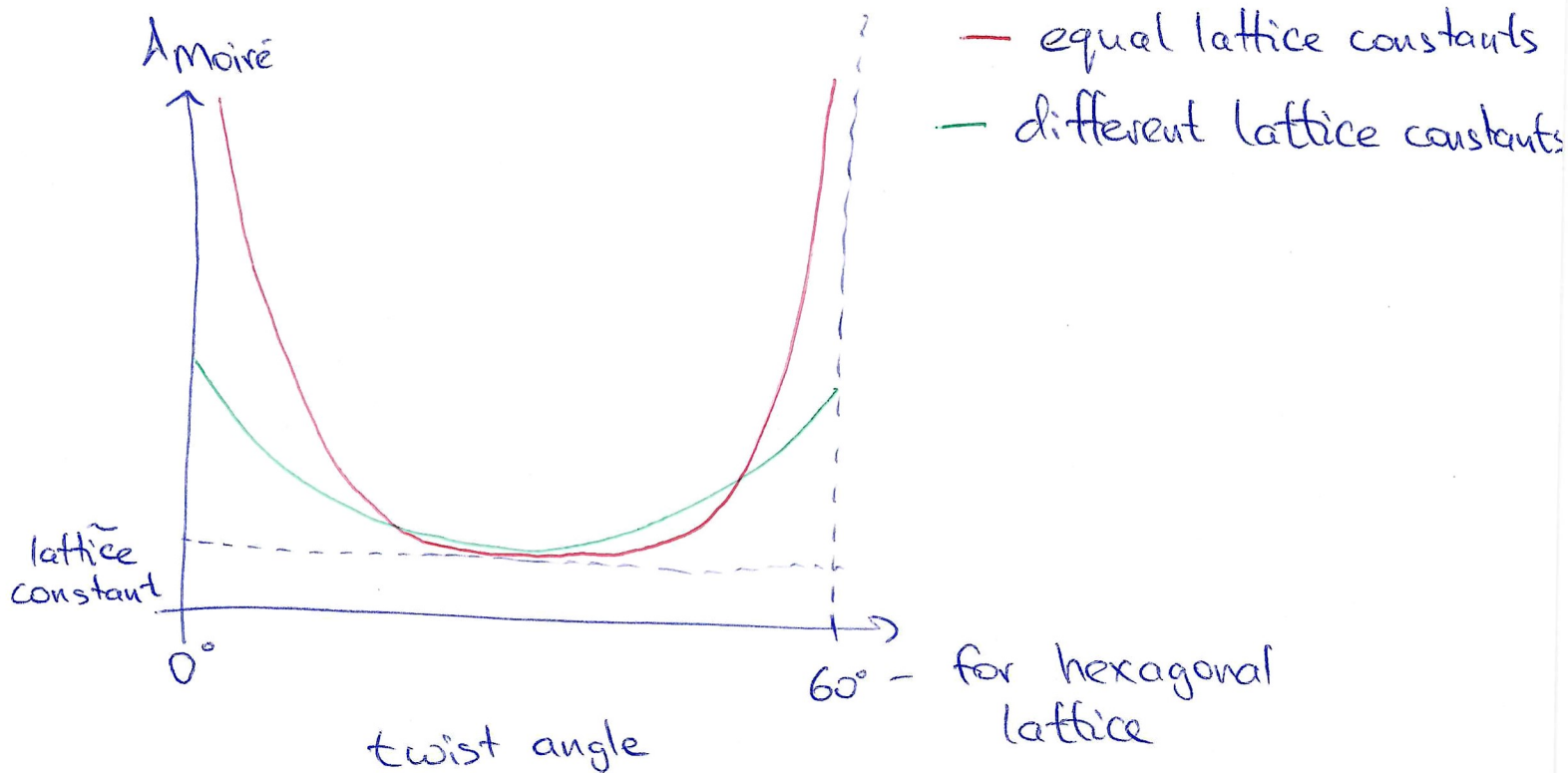
→ appears when two inequivalent lattices are overlaid



- o lattice constant mismatch
- o angle mismatch



↳ Moiré superlattice period:



[→ however, at very small angles atomic reconstruction can occur (atoms shift slightly into energetically more favorable positions)
↳ different patterns & domains emerge]

⇒ rich, contemporary field of 2D materials physics studying electronic, optical & vibrational excitations in artificial lattices