

- Crude approximation $U(t+\delta t, t) \approx e^{-iH(t)\delta t}$ can be improved by higher order commutator-free expansion (see, e.g. A. Alvermann, H. Fehske, Journal of Computational Physics 230, 5930 (2011), arXiv: 1102.5071)

→ Result up to fourth order:

$$U(t+\delta t, t) = e^{\frac{1}{2}A_1 + \frac{1}{3}A_2} e^{\frac{1}{2}A_1 - \frac{1}{3}A_2} + \mathcal{O}(\delta t^5)$$

with $A_1 = -i \int_t^{t+\delta t} H(\tilde{c}) d\tilde{c}$ and $A_2 = -3i \int_t^{t+\delta t} (2\frac{\tilde{c}}{\delta t} - 1) H(\tilde{c}) d\tilde{c}$.

→ Integrals are performed numerically with a single Runge-Kutta 4 step to yield:

$$A_1 \approx \hat{1} + \mathcal{O}(\delta t^5) - i \frac{\delta t}{6} \left(H(t) + H(t+\delta t) + 4H(t + \frac{\delta t}{2}) \right)$$

$$A_2 \approx -i \frac{\delta t}{2} \left[\left(2 \frac{t}{\delta t} - 1 \right) H(t) + \left(2 \frac{t}{\delta t} + 1 \right) H(t + \delta t) + 8 \frac{t}{\delta t} H\left(t + \frac{\delta t}{2}\right) \right].$$

↳ Practical recipe:

(i) Compute $K_{n+1} \left(i \left(\frac{A_1}{2} - \frac{A_2}{3} \right), |\Psi(t)\rangle \right)$ with basis Q_- and tridiagonal form T_- .

(ii) Update $|\tilde{\Psi}\rangle = Q_- e^{-iT_-} Q_-^\dagger |\Psi(t)\rangle$

(iii) Compute $K_{n+1} \left(i \left(\frac{A_1}{2} + \frac{A_2}{3} \right), |\tilde{\Psi}\rangle \right) \tilde{y}_0(t)$ with basis Q_+ and T_+ .

(iv) Update $|\Psi(t + \delta t)\rangle = Q_+ e^{-iT_+} Q_+^\dagger \tilde{\Psi} \approx \tilde{y}_0$

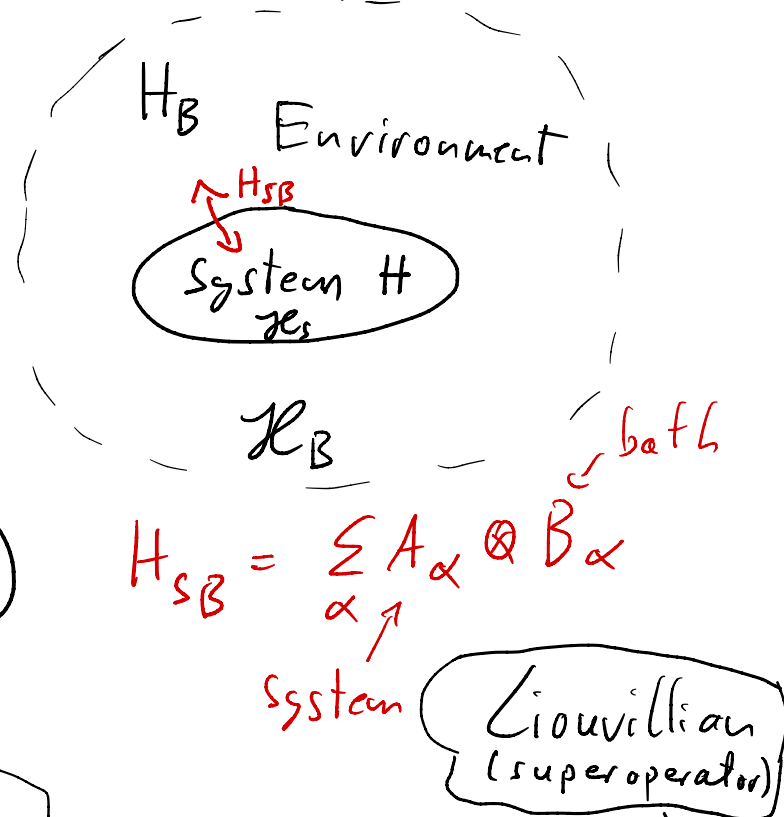
Continue with next time-step.

B.5: Quantum Master Equations and Dissipative Dynamics

- Motivation: System coupled to its environment (bath).

↳ Key ingredients:

- weak coupling (Born)
- bath without memory (Markov)



↳ Lindblad master equation:

$\frac{d}{dt} \rightarrow$

for pure states $\rho = |\psi\rangle\langle\psi|$

$$\dot{\rho} = -i[H, \rho] + \sum_{\lambda} \left(L_{\lambda} \rho L_{\lambda}^{\dagger} - \frac{1}{2} \{ L_{\lambda}^{\dagger} L_{\lambda}, \rho \} \right) = \mathcal{L}(\rho)$$

Operator on $\mathcal{H}_S$

- Vectorize Lindblad master equation as

$$|\dot{\psi}_\rho\rangle = \mathcal{L} |\psi_\rho\rangle, \text{ where } |\psi_\rho\rangle = \sum_{i,j} \rho_{ij} \underbrace{|i\rangle}_{\in \mathcal{H}_S} \otimes \underbrace{|j\rangle}_{\in \mathcal{H}_B = \mathcal{H}_S}$$

for $\rho = \sum_{i,j} \rho_{ij} |i\rangle\langle j|$

is the vectorized state on which $\mathcal{L}$ acts as a linear operator.

In coordinate representation, we have:

$$\langle l, m | \dot{\Psi}_S \rangle = \dot{\rho}_{lm} = -i(H_{en} \rho_{nm} - \rho_{en} H_{nm}) + \sum_n \left[L_n^{en} \rho_{np} L_n^{+pm} - \frac{1}{2} \left((L_n^\dagger L_n)_{en} \rho_{nm} + \rho_{en} (L_n^\dagger L_n)_{nm} \right) \right].$$

In vectorized form, this yields

$$|\dot{\Psi}_S\rangle = \mathcal{L} |\Psi_S\rangle \text{ with } \mathcal{L} = -i \left(\overset{\mathcal{H}_S}{H} \otimes \mathbb{1} - \mathbb{1} \otimes \overset{\mathcal{H}_A}{H}^T \right) + \sum_n \left[L_n \otimes L_n^\dagger - \frac{1}{2} \left((L_n^\dagger L_n) \otimes \mathbb{1} + \mathbb{1} \otimes (L_n^T L_n^*) \right) \right].$$

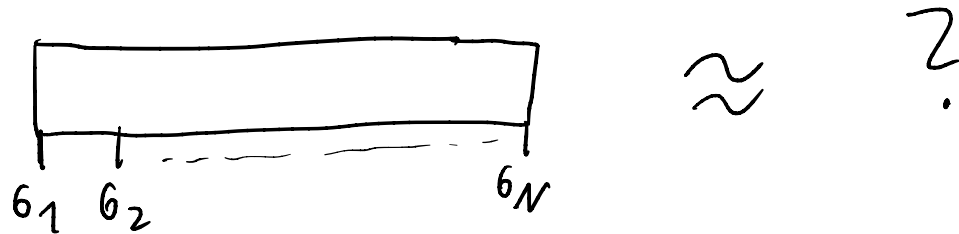
↳ Price to pay:

- Linear operator $\mathcal{L}$ non-Hermitian (Lanczos $\rightarrow$ Arnoldi)
- $\dim(\mathcal{H}_S) \rightarrow \dim(\mathcal{H}_S \otimes \mathcal{H}_A) = (\dim(\mathcal{H}_S))^2$
(effective doubling in system size)

↳ Remedy in part D: Stochastic sampling of dissipative "jumps" (action of L_i 's) allows for algorithm based on pure states in $\mathcal{H}_S$ (quantum trajectories), and thus brings us back to $\dim(\mathcal{H}_S)$.

C.O : The variational principle

- Basic idea: Tame exponential complexity of $|\psi\rangle$ by making a clever variational ansatz



- ↳ Key challenges:
- (i) Identify set of variational states that is unbiased and contains states complex enough that they are close to the true solution
 - (ii) Devise constructive algorithm that optimizes the variational parameters.

