

• Bringing an MPS into canonical forms:

- "left canonization" of generic MPS:

$$|\varphi\rangle = \sum_{\{b\}} A_{1a_1}^{b_1} A_{a_1a_2}^{b_2} \dots A_{a_{L-1}1}^{b_L} |b_1, \dots, b_L\rangle =$$

$$U_{b_1, s_1} = L_{1a_1}^{b_1}$$

$$= \sum_{\{b\}} L_{1s_1}^{b_1} \underbrace{\left(\sum_{a_1} S_{s_1 s_1} V_{s_1 a_1}^\dagger A_{a_1 a_2}^{b_2} \right)}_{\substack{\sim b_2 \\ =: \tilde{A}_{s_1 a_2}^{b_2}}} A_{a_2 a_3}^{b_3} \dots A_{a_{L-1} 1}^{b_L} |b_1, \dots, b_L\rangle$$

$$\begin{array}{c} \text{SVD} \\ \downarrow \\ A_{1a_1}^{b_1} = U_{b_1, s_1} S_{s_1 s_1} V_{s_1 a_1}^\dagger \end{array}$$

SVD

$$\tilde{A}_{s_1 a_2}^{b_2} \rightarrow \sum_{\{b\}} L_{1s_1}^{b_1} L_{s_1 s_2}^{b_2} \tilde{A}_{s_2 a_3}^{b_3} \dots |b_1, \dots, b_L\rangle = \dots$$

$$\begin{aligned}
 \dots &= \sum_{\{b\}} L_{1s_1}^{b_1} L_{s_1 s_2}^{b_2} \dots L_{s_{L-1} s_L}^{b_L} \underbrace{S_{11} V_{11}^+}_{\text{scalar} \sim \text{Norm of } |\psi\rangle} |b_1, \dots, b_L\rangle
 \end{aligned}$$

$\hookrightarrow$ Crucial feature of left-canonical MPS
 (by def. of SVD):

$$\sum_{G_j} L_{G_j}^{b_j+} L_{G_j}^{b_j} = \mathbb{1} \quad \forall_j$$

(equivalent to $U^+ U = \mathbb{1}$)
 in SVD

- "right canonization" of generic MPS: Analogous
 to left-canonical case, but using $A^b = U S V^+$,
 where V^+ is used to define R^{b_j} such that

$$|\psi\rangle = \sum_{\{b\}} \underbrace{U_{11} S_{11}}_{\text{scalar} \sim \text{norm } |\psi\rangle} R_{1s_1}^{b_1} R_{s_1 s_2}^{b_2} \dots R_{s_{L-1}, 1}^{b_L} |b_1, \dots, b_L\rangle$$

with

$$\sum_{b_j} R^{b_j} R^{b_j \dagger} = \mathbb{1} \quad \forall_j$$

- Mixed canonical MPS: Left-canonical up to site l , right-canonical from $l+1$ to L :

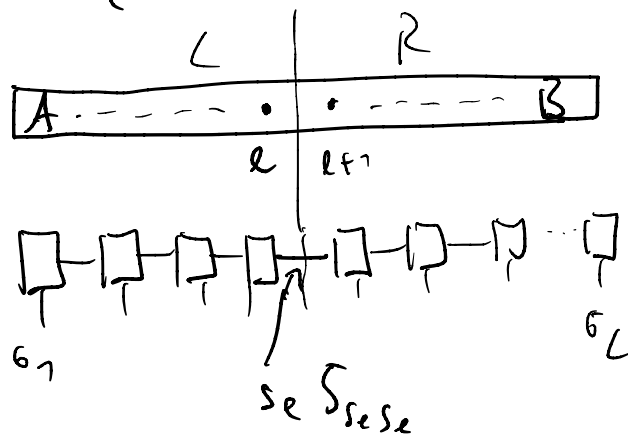
$$|\psi\rangle = \sum_{\{b\}} L_{1s_1}^{b_1} L_{s_1 s_2}^{b_2} \dots L_{s_{l-1} s_l}^{b_l} S_{s_l s_l} R_{s_l s_{l+1}}^{b_{l+1}} \dots R_{s_{L-1} s_L}^{b_L} |b_1, \dots, b_L\rangle$$

$\hookrightarrow$ With $|S_l^A\rangle = \sum_{b_1, \dots, b_l} \left(L^{b_1} \dots L^{b_l} \right)_{1s_l} |b_1, \dots, b_l\rangle$ and

$$|S_l^B\rangle = \sum_{b_{l+1}, \dots, b_L} \left(R^{b_{l+1}} \dots R^{b_L} \right)_{s_{l+1}} |b_{l+1}, \dots, b_L\rangle,$$

we get $|\psi\rangle = \sum_{s_e} S_{s_e} |s_e^A\rangle \otimes |s_e^B\rangle$

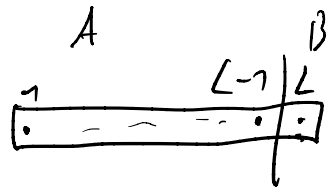
Graphically:



$\rightarrow$ mixed canonical form makes Schmidt decomposition manifest, since $\sum_{b_i} L^{b_i\dagger} L^{b_i} = \mathbb{1}$, $i=1, \dots, l$ and $\sum_{b_i} R^{b_i} R^{b_i\dagger} = \mathbb{1}$, $i=l+1, \dots, L$ ensure orthonormality of the $\{|s_e^A\rangle\}$ and $\{|s_e^B\rangle\}$ basis sets.

- Compress MPS by SVD (from bond dim. χ to $\chi < \chi$)
- Starting point: Left-canonical MPS $|\psi^{(L)}\rangle = \sum_{\{b\}} L^{b_1} \dots L^{b_L} |b_1 \dots b_L\rangle$
- Basic procedure: Right canonicalize $|\psi^{(L)}\rangle$ site by site, retaining only the χ largest singular values at every SVD step:

$$|\psi^{(L-1)}\rangle = \sum_{\{b\}} L^{b_1} \dots L^{b_{L-1}} U S R^{b_L} |b_1 \dots b_L\rangle$$



$\hookrightarrow$ Read off Schmidt-decomp.: $|S_{L-1}^B\rangle = \sum_{b_L} R_{S_{L-1}, b_L}^{b_L} |b_L\rangle$

$$|S_{L-1}^A\rangle = \sum_{b_1, \dots, b_{L-1}} (L^{b_1} \dots L^{b_{L-1}} U)_{1, S_{L-1}} |b_1, \dots, b_{L-1}\rangle \quad \text{i.e.}$$

$$|\psi^{(L-1)}\rangle = \sum_{S_{L-1}} \sum_{s_{L-1}} |S_{L-1}^A\rangle \otimes |S_{L-1}^B\rangle \quad (\text{keep only } s_{L-1} = 0, \dots, \chi)$$

Continue from right to left,
at every step, truncate matrices to only keep the
 χ largest singular values

until $|\tilde{\varphi}\rangle = \sum_{\{b\}} \tilde{R}^{b_1} \dots \tilde{R}^{b_L} |b_1, \dots, b_L\rangle$ is reached.

$\leadsto |\tilde{\varphi}\rangle$ is the compressed MPS of bond dim. χ .

— Problem with SVD compression: Truncation procedure has clear direction that may lead to artefacts for "lossy compressions" (significant singular values being cut).

$\leadsto$ Cure: More elaborate iterative optimization algorithms for compression (see U. Schollwöck 2011 Review for details)

• Matrix product operators (MPOs)

- Basic idea: By analogy to MPS form

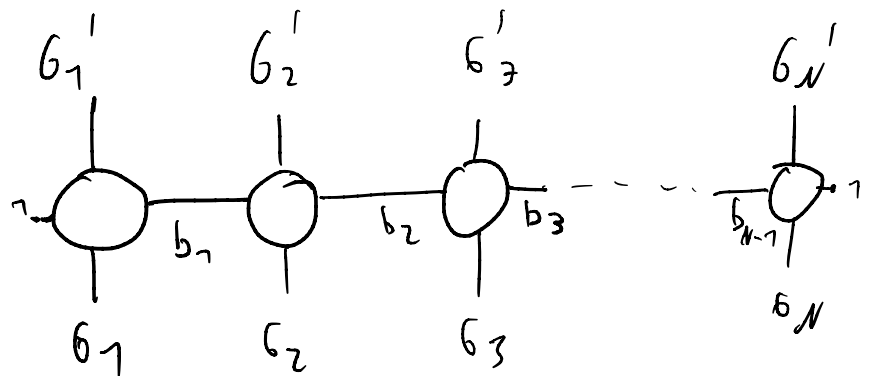
$$\langle b_1, \dots, b_N | \Psi \rangle = A^{b_1} \dots A^{b_N},$$

represent $\langle b_1, \dots, b_N | O | b'_1, \dots, b'_N \rangle = W^{b_1 b'_1} \dots W^{b_N b'_N}$

with matrices $W_{b_{j-1} b_j}^{b'_j b'_j}$, resulting in the form

$$O = \sum_{\{b, b'\}} W_{1 b_1}^{b'_1 b'_1} W_{b_1 b_2}^{b'_2 b'_2} \dots W_{b_{N-1} b_N}^{b'_N b'_N} \langle b'_1, \dots, b'_N | \langle b_1, \dots, b_N |$$

or graphically

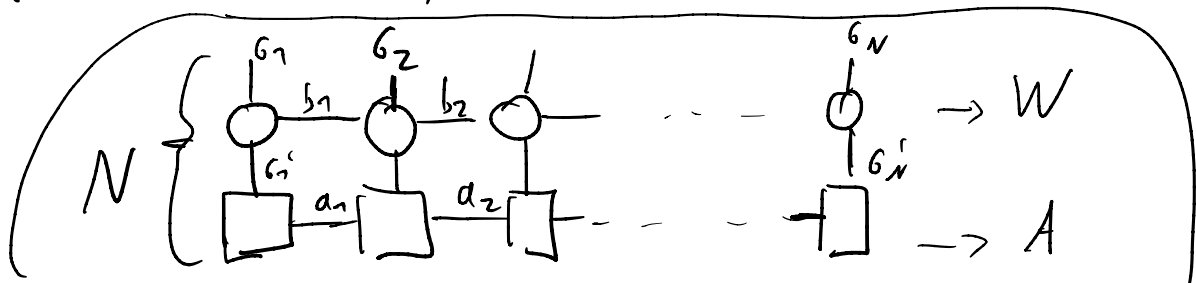


↳ At the expense of a possibly exponentially big bond dimension (size of W matrices for fix physical indices), any operator can be brought into MPO form by successive SVD.

↳ Applying an MPO to an MPS :

$$O|\psi\rangle = \sum_{\{b, b'\}} (W^{b_1 b_1'} \dots W^{b_N b_N'}) (A^{b_1'} \dots A^{b_N'}) |b_1, \dots, b_N\rangle$$

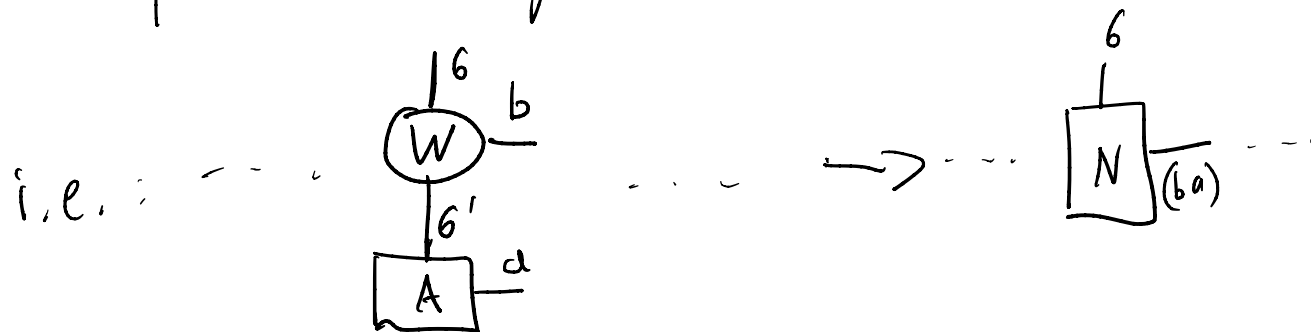
$$= \sum_{\substack{a_1, \dots, a_{N-1} \\ b_1, \dots, b_{N-1}}} \sum_{\{b, b'\}} \left(W_{b_1 b_1'}^{b_1 b_1'} A_{a_1 a_1}^{b_1'} \right) \left(W_{b_1 b_2}^{b_2 b_2'} A_{a_1 a_2}^{b_2'} \right) \dots \left(W_{b_{N-1} b_N}^{b_N b_N'} A_{a_{N-1} a_N}^{b_N'} \right) |b_1, \dots, b_N\rangle$$



$$= \sum_{b_1 a_1, \dots, b_{N-1} a_{N-1}} \sum_{\{b\}} N_{1(b_1 a_1)}^{b_1} N_{(b_1 a_1)(b_2 a_2)}^{b_2} \dots N_{(b_{N-1} a_{N-1})1}^{b_N} |b_1, \dots, b_N\rangle$$

i.e. $|\psi\rangle$ is MPS with $N_{(b_{i-1} a_{i-1})(b_i a_i)}^{b_i} = \sum_{b'_i} W_{b_{i-1} b_i}^{b_i b'_i} A_{a_i a_i}^{b'_i}$.

$\leadsto$ The bond dimension of $|\psi\rangle$ is the product of the bond-dim. of MPS " $|\psi\rangle$ " and MPO " 0 ", and the N matrices are tensor products of A and W matrices: $N^{b_i} = \sum_{b'_i} W^{b_i b'_i} \otimes A^{b'_i}$.



In generic algorithms such as tMPS, MPOs with small individual bond dim. are iteratively applied to an MPS, and the resulting MPS is compressed after every step.